



Introduction

A very important class of special functions that appear in many calculations of multiloop Feynman integrals are the so-called multiple polylogarithms (MPLs), defined as

$$G(a_1, a_2, \dots, a_n; x) = \int_0^x \frac{dt}{t - a_1} G(a_2, \dots, a_n; t). \quad (1)$$

The vector $\vec{a} = (a_1, \dots, a_n)$ constitutes the *indices* of the MPLs and is also called the *vector of singularities*. There are specific cases where MPLs can be written as more basic functions, such as logarithms, classical polylogarithms $\text{Li}_n(x)$ and Nielsen polylogarithms $S_{n,p}(x)$

$$G(\vec{0}_n; x) = \frac{1}{n!} \log^{(n)}(x) \quad (2)$$

$$G(\vec{a}_n; x) = \frac{1}{n!} \log^{(n)}\left(1 - \frac{x}{a}\right) \quad (3)$$

$$G(\vec{0}_{n-1}, a; x) = -\text{Li}_n\left(\frac{x}{a}\right) \quad (4)$$

$$G(\vec{0}_n, \vec{a}_p; z) = (-1)^p S_{n,p}\left(\frac{x}{a}\right). \quad (5)$$

In the following we will denote with $f(a; x) = \frac{1}{x-a}$ the integration kernel. We will also denote as $G(\vec{a}_w; x)$ a weight- w MPL, i.e. a MPL whose vector of singularities has w entries. Unless stated otherwise we will consider $0 < x < 1$.

Argument transformations, derivatives and analytic continuation

1. Show that

$$G(0, \vec{a}_{w-1}; 1-x) = G(0, \vec{a}_{w-1}; 1) + \int_0^x \frac{dy}{y-1} G(\vec{a}_{w-1}; 1-y) \quad (6)$$

and then use the transformation to express $G(0, 0, 1, 1; 1-x)$ in terms of MPLs with argument x .

2. Show that

$$G\left(\vec{a}; \frac{1}{x}\right) = G(a_1, \vec{a}_{w-1}; 1) - \int_1^x \frac{d}{t^2} f\left(a_1; \frac{1}{t}\right) G\left(\vec{a}_{w-1}; \frac{1}{t}\right). \quad (7)$$

Argue that we can in general assume that $a_1 \neq 1$. In this case $G(a_1, \vec{a}_{w-1}; 1)$ is a finite constant¹. For this particular transformation, one must also pay attention to the branch points of the MPLs. Using $\log(-x) = \log(x) - i\pi$ for the analytic continuation of the logarithm, apply the above transformation to $G\left(0, -1, 1; \frac{1}{x}\right)$.

¹When we say constant, we mean in the variable x . the vector \vec{a} could in general depend on other variables.

3. Show that

$$G\left(0, \vec{a}_{w-1}; \frac{1-x}{1+x}\right) = G(0, \vec{a}_{w-1}; 1) - \int_0^x dt \left(\frac{1}{t-1} + \frac{1}{t+1} \right) G\left(\vec{a}_{w-1}; \frac{1-t}{1+t}\right) \quad (8)$$

$$G\left(-1, \vec{a}_{w-1}; \frac{1-x}{1+x}\right) = G(-1, \vec{a}_{w-1}; 1) - \int_0^x \frac{dt}{t+1} G\left(\vec{a}_{w-1}; \frac{1-t}{1+t}\right) \quad (9)$$

and use these transformations to express $G(-1, -1, 1; \frac{1-x}{1+x})$ in terms of MPLs with argument x .

4. Assuming that all arguments of the MPLs are generic, prove that

$$dG(a_1; x) = d\log\left(\frac{x-a_1}{-a_1}\right) \quad (10)$$

$$dG(a_1, a_2; x) = G(a_2; x)d\log\left(\frac{x-a_1}{a_2-a_1}\right) + G(a_1; x)d\log\left(\frac{a_1-a_2}{-a_2}\right). \quad (11)$$

What is the total differential of a weight three MPL, i.e. $dG(a_1, a_2, a_3; x)$? Starting from the explicit cases above, can you write down a general formula for the total differential of a MPL of arbitrary weight?

5. Consider the solution of the canonical massive one loop bubble in $D = 4 - 2\epsilon$ dimensions up to weight two,

$$I(x, \epsilon) = \epsilon G(0; x) + \epsilon^2 [G(0, 0; x) - 2G(-1, 0; x)]. \quad (12)$$

To write the analytic result, we used the Landau variable defined in the euclidean region

$$p^2 = m^2 \frac{(1-x)^2}{x}, \quad \text{where } p^2 \geq 0 \rightarrow 0 \leq x \leq 1.$$

We would like now to continue the result *above threshold*, i.e. for $s = -p^2 \leq 4m^2$. To do that, we continue $x \rightarrow -y \pm i\delta$, where δ is a small, positive imaginary part, such that

$$s = m^2 \frac{(1+y)^2}{y}, \quad \text{where } s \geq 4m^2 \rightarrow 0 \leq y \leq 1.$$

- (a) Determine the sign of the imaginary part associated to y , by using the fact that s must take a positive imaginary part when above threshold ($s + i\delta$).
- (b) With the proper sign for the imaginary part, perform the analytic continuation for the one-loop bubble and prove that

$$I(y, \epsilon) = \epsilon [G(0; y) - i\pi] + \epsilon^2 \left[-i\pi G(0; y) + 2i\pi G(1; y) + G(0, 0; y) - 2G(1, 0; y) - \frac{\pi^2}{2} \right]. \quad (13)$$