Scattering Amplitudes in QFT WS 2023/24

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Sheet 04: Four-gluon amplitude squared and colour decomposition at higher points and loop level

Exercise 1 - Four gluon amplitude squared and summed over colours and polarisations

In the last exercise sheet, we derived the colour-decomposed tree-level four-gluon amplitude

$$\mathcal{M}_{4g} = 4 \sum_{\sigma \in S_3} \operatorname{tr} \left(T^{a_1} T^{a_{\sigma(2)}} T^{a_{\sigma(3)}} T^{a_{\sigma(4)}} \right) \,\mathcal{A} \left[1, \sigma(2), \sigma(3), \sigma(4) \right] \,, \tag{1}$$

where the MHV amplitudes read

$$\mathcal{A}\left[1^{+}, 2^{+}, 3^{-}, 4^{-}\right] = g^{2} \frac{\langle 12 \rangle^{4}}{\langle 12 \rangle \langle 23 \rangle \langle 34 \rangle \langle 41 \rangle},$$

$$\mathcal{A}\left[1^{+}, 2^{-}, 3^{+}, 4^{-}\right] = g^{2} \frac{\langle 13 \rangle^{4}}{\langle 12 \rangle \langle 23 \rangle \langle 34 \rangle \langle 41 \rangle},$$

$$\mathcal{A}\left[1^{+}, 2^{-}, 3^{-}, 4^{+}\right] = g^{2} \frac{\langle 14 \rangle^{4}}{\langle 12 \rangle \langle 23 \rangle \langle 34 \rangle \langle 41 \rangle}.$$

(2)

In the lecture, it was discussed that we can extend $SU(N) \rightarrow U(N)$ for tree-level gluon amplitudes since they do not receive contributions from fermions. This allows us to use $T_{ij}^a T_{kl}^a = \delta_{il} \delta_{kj}/2$ when manipulating strings of colour generators.

1. First, show that the following identities hold:

$$\operatorname{tr}(T^{a}A)\left[\operatorname{tr}(T^{a}B)\right]^{\dagger} = \frac{1}{2}\operatorname{tr}(AB^{\dagger}),$$

$$\operatorname{tr}(T^{a}AT^{a}B) = \frac{1}{2}\operatorname{tr}(A)\operatorname{tr}(B).$$
(3)

2. Using Eq.(3), show that squared traces in $\sum_{col} |\mathcal{M}_{4g}|^2$ yield

$$\operatorname{tr}\left(T^{a_{1}}T^{a_{\sigma(2)}}T^{a_{\sigma(3)}}T^{a_{\sigma(4)}}\right)\left[\operatorname{tr}\left(T^{a_{1}}T^{a_{\sigma(2)}}T^{a_{\sigma(3)}}T^{a_{\sigma(4)}}\right)\right]^{\dagger} = \frac{N^{4}}{16}.$$
(4)

3. Using Eq.(3), show that interference terms of traces in $\sum_{col} |\mathcal{M}_{4g}|^2$ yield

$$\operatorname{tr}\left(T^{a_{1}}T^{a_{\sigma(2)}}T^{a_{\sigma(3)}}T^{a_{\sigma(4)}}\right)\left[\operatorname{tr}\left(T^{a_{1}}T^{a_{\sigma'(2)}}T^{a_{\sigma'(3)}}T^{a_{\sigma'(4)}}\right)\right]^{\dagger} = \frac{N^{2}}{16},$$
(5)

for $\sigma \neq \sigma'$ different permutations.

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4. Using the results for color traces in Eq.(4) and Eq.(5), we can write

$$\sum_{\text{col}} |\mathcal{M}_{4g}| = N^4 \sum_{\sigma \in S_3} |\mathcal{A}[1, \sigma(2), \sigma(3), \sigma(4)]|^2 + N^2 \sum_{\substack{\sigma, \sigma' \in S_3 \\ \sigma \neq \sigma'}} \mathcal{A}[1, \sigma(2), \sigma(3), \sigma(4)] (\mathcal{A}[1, \sigma'(2), \sigma'(3), \sigma'(4)])^{\dagger}.$$
(6)

Using the relation between the colour-ordered amplitudes,

- (a) Cyclicity: $\mathcal{A}[a, b, c, d] = \mathcal{A}[d, a, b, c] = \mathcal{A}[c, d, a, b] = \mathcal{A}[b, c, d, a]$, (7)
- (b) Reflection symmetry: $\mathcal{A}[a, b, c, d] = \mathcal{A}[d, c, b, a]$, (8)
- (c) Photon decoupling: $\mathcal{A}[a, b, c, d] + \mathcal{A}[b, a, c, d] + \mathcal{A}[b, c, a, d] = 0.$ (9)

show that we can write

$$\sum_{\text{col}} |M_{4g}| = 2N^2(N^2 - 1) \left\{ |\mathcal{A}[1, 2, 3, 4]|^2 + |\mathcal{A}[1, 2, 4, 3]|^2 + |\mathcal{A}[1, 3, 2, 4]|^2 \right\}.$$
 (10)

5. Compute the sum over helicities for the first partial amplitude in Eq.(10) using the results in Eq.(2). You should find

$$\sum_{\text{pol}} |\mathcal{A}[1,2,3,4]|^2 = 2g^4 \left[\frac{s_{12}^2}{s_{23}^2} + \frac{s_{13}^4}{s_{12}^2 s_{23}^2} + \frac{s_{23}^2}{s_{12}^2} \right].$$
(11)

6. Use the result in Eq.(11) to compute the final result

$$\sum_{\text{pol,col}} |M_{4g}(1,2,3,4)| = 2g^4 N^2 (N^2 - 1) \times \frac{(s_{12}^2 + s_{13}^2 + s_{14}^2)^3}{s_{12}^2 s_{13}^2 s_{14}^2}.$$
 (12)

Exercise 2 - Colour ordering in QCD at higher point and higher loops

1. Argue that tree-level n-gluon amplitudes admit for any n a colour decomposition of the form

$$\mathcal{M}_{ng}^{\text{tree}} = \sum_{\sigma \in S_{n-1}} \text{tr} \left(T^{a_1} T^{a_{\sigma(2)}} \dots T^{a_{\sigma(n)}} \right) \mathcal{A} \left[1, \sigma(2), \dots \sigma(n) \right] \,. \tag{13}$$

- 2. Argue that extending $SU(N) \to U(N)$ is no longer possible at loop level.
- 3. Argue that the colour basis must be enlarged at loop-level. Give an explicit example of a Feynman diagram leading to this enlargement at n = 4 and derive the type of new colour basis elements produced by it.