

# Relativity, Particles, Fields SS 2017

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<https://www.t75.ph.tum.de/teaching/ss17-relativity-particles-fields/>

Sheet 5: Canonical quantization (30.5.2017)



## 1 The charge of a state

Consider a complex scalar field operator  $\psi$ . Assume that  $|\alpha\rangle$  is an eigenstate of the charge operator  $Q$  with eigenvalue  $q$ . Prove that the state obtained by applying the operator  $\psi^\dagger$ , i.e.  $\psi^\dagger|\alpha\rangle$ , is also an eigenstate of  $Q$ . What is the corresponding eigenvalue?

## 2 Fluctuations in two dimensions

Consider a massless real scalar field  $\phi$  in two-dimensional spacetime.

a) Show that given two space-time points  $x$  and  $y$ , the fluctuations take the following form

$$\langle 0|\phi(x)\phi(y)|0\rangle = \frac{1}{4\pi} \int_0^\infty \frac{dk}{k} (e^{ik(\tau-r)} + e^{ik(\tau+r)}) \quad (1)$$

where  $\tau \equiv y^0 - x^0$  and  $r \equiv y^1 - x^1$ .

b) Show that Eq. (1) can be written in the explicit form

$$\langle 0|\phi(x)\phi(y)|0\rangle = -\frac{1}{4\pi} \log \frac{(x-y)^2}{\mu^2}, \quad (2)$$

where  $\mu$  is an integration constant which has the dimension of length.

*Hint:* a convenient way to proceed is to differentiate the RHS of Eq. (1) with respect to  $\tau$ . The resulting integral can be regularized by multiplying the integrand times  $e^{-\epsilon k}$  with  $\epsilon > 0$ , and eventually taking the limit  $\epsilon \rightarrow 0$ .

## 3 Propagator in position space

Consider a scalar field  $\phi$  with mass  $m$ , whose Feynman propagator is

$$\Delta_F(x) = \langle 0|T\phi(x)\phi(0)|0\rangle. \quad (3)$$

a) Show that  $\langle 0|\phi(x)|\mathbf{k}\rangle = e^{-i\omega_k t + i\mathbf{k}\cdot\mathbf{r}}$ , where  $|\mathbf{k}\rangle$  is a one-particle state and  $x^\mu = (t, \vec{r})$ . Then use this fact to prove that the propagator can be written in the integral form

$$\Delta_F(x) = \frac{1}{4\pi^2} \int_0^\infty dk \frac{k^2 \sin kr}{\omega_k kr} e^{-i\omega_k |t|}, \quad (4)$$

with  $r = |\vec{r}|$ .

b) By Lorentz invariance,  $\Delta_F$  can depend only on  $x^2 = t^2 - r^2$ . Assuming a space-like separation  $x^2 < 0$ , set  $t = 0$  and plot: 1)  $4\pi^2 r^2 \Delta_F(r)$  in the range  $0 < r < 2$ , for  $m = 1$  and  $m = 3$ . Observe the  $r \rightarrow 0$  limit. Does it depend on  $m$ ? 2)  $4\sqrt{2} \pi^{3/2} \frac{r^{3/2}}{\sqrt{m}} \Delta_F(r)$  again in the range  $0 < r < 2$ , for  $m = 2$  and  $m = 3$ . At

large  $r$ , this tends to an elementary function of  $mr$ ,  $f(mr)$ . Determine  $f$  and conclude that the typical correlation length is roughly  $\Delta r \sim 1/m$ .

c) Now set  $m = 0$ , and prove that in this case the propagator takes the general form

$$\Delta_F(x)_{m=0} = -\frac{i}{4\pi}\delta(x^2) - \frac{1}{4\pi^2}\text{P}\frac{1}{x^2}, \quad (5)$$

where P denotes the principal value prescription. To this end, recall that

$$\frac{1}{x \pm i\epsilon} = \text{P}\frac{1}{x} \mp i\pi\delta(x) \quad (6)$$

where  $\epsilon > 0$ .