Übungsblatt Nr. 12

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Mechanik (Theoretische Physik 1)

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Aufgabe 1: Noether charge with transformations of time

As seen in the lectures, if the action is invariant under an infinitesimal transformation of the generalized coordinates $q_i(t) \rightarrow q'_i(t) = q_i(t) + \epsilon \delta Q_i[q_i(t)]$, such that

$$dt\left(\mathcal{L}\left[t,q_{i}'[t],\frac{dq_{i}'}{dt}[t]\right]-\epsilon\frac{dR}{dt}\right)=dt\mathcal{L}\left[t,q_{i}[t],\frac{dq_{i}}{dt}[t]\right],$$
(1)

then there is a conserved Noether charge

$$Q = \sum_{i} \frac{\partial \mathcal{L}}{\partial \dot{q}_{i}} \delta Q_{i} - R.$$
⁽²⁾

Consider now transformations of both the generalized coordinates and the time variable,

$$q_i(t) \to q'_i(t') = q_i(t) + \epsilon \delta Q[q_i(t)],$$

$$t \to t' = t + \epsilon \delta T(t),$$

which are assumed to leave the action invariant.

1. Show that one has (up to higher order corrections in ϵ)

$$dt'\mathcal{L}\left[t',q_i'[t'],\frac{dq_i'}{dt'}[t']\right] = dt\left(1+\epsilon\frac{d\delta T}{dt}\right)\mathcal{L}\left[t+\epsilon\delta T,q_i'[t'],\frac{dq_i'}{dt}[t']-\epsilon\frac{d\delta T}{dt}\frac{dq_i}{dt}[t]\right] = dt\left(\mathcal{L}\left[t,q_i'[t'],\frac{dq_i'}{dt}[t']\right]-\epsilon\frac{d}{dt}\left[\left(\sum_i\frac{\partial\mathcal{L}}{\partial\dot{q}_i}\dot{q}_i-\mathcal{L}\right)\delta T\right]\right),$$

which, according to equations (1) and (2), leads to the conservation of the charge

$$Q = \sum_{i} \frac{\partial \mathcal{L}}{\partial \dot{q}_{i}} \delta Q_{i} - \left(\sum_{i} \frac{\partial \mathcal{L}}{\partial \dot{q}_{i}} \dot{q}_{i} - \mathcal{L}\right) \delta T.$$

[Hints. Use the identity

$$\frac{\partial \mathcal{L}}{\partial t} = \frac{d\mathcal{L}}{dt} - \sum_{i} \left(\dot{q}_i \frac{\partial \mathcal{L}}{\partial q_i} + \ddot{q}_i \frac{\partial \mathcal{L}}{\partial \dot{q}_i} \right),$$

as well as the Euler-Lagrange equations.]

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2. Consider the Lagrangian of a particle in a central potential decaying as $1/r^2$:

$$\mathcal{L} = \frac{1}{2}m\dot{r}^2 - \frac{\alpha}{r^2}.$$

Show that the action is invariant under transformations

$$r[t] \rightarrow r'[t'] = (1+\epsilon)r[t],$$

$$t \rightarrow t' = (1+\epsilon)^2 t,$$

and compute the associated Noether charge in terms of energy, radial coordiante and its time derivative.

Hint: Remember that when the time coordinate is affected by the transformation, the measure of integration that appears in the action should change as well, namely:

$$dt \rightarrow \Big| \frac{dt}{dt'} \Big| dt'$$

Aufgabe 2: The Steiner theorem

Consider a frame K centered in the center of mass of a rigid body of mass M and fixed to the body. In this frame, the tensor of inertia of the rigid body is $J_{\mu\nu}$. If we consider another frame K' which is parallel to K but shifted by a constant vector **a**, the tensor of inertia in this frame will be $J'_{\mu\nu}$. Show that the two tensors are related by the so called Steiner theorem for parallel axes:

$$J'_{\mu\nu} = J_{\mu\nu} + M[|\mathbf{a}|^2 \delta_{\mu\nu} - a_\mu a_\nu]$$

Keeping in mind the latter relation, complete the following tasks:

- 1. Compute the tensor of inertia of a sphere of radius R and homogeneous mass M in the frame of its center of mass.
- 2. Consider a system made of two spheres with equal mass M/2 and same radius R, welded together in their point of contact. Compute the tensor of inertia in centered in the latter.
- 3. Compute the tensor of inertia of a regular cone of height h, and base radius R in the frame centered in its center of mass. (The position of the center of mass must be calculated.) Hint1: In the frame centered in the vertex of the cone and with one axis along the symmetry axis, the tensor of inertia is already diagonal. Compute I_i , i = 1, 2, 3 in that frame and then use the Steiner theorem to shift them to the center of mass. Hint2: Cylindrical coordinates might help.
- 4. Consider a homogeneous cuboid of mass M and dimensions $a_1 = a_2$, a_3 . Starting from frame K shown in figure, the tensor of inertia is diagonal. Compute it in that frame, and then compute it in the primed frame K' in which the x'_3 axis is pointing in the direction of the main diagonal of the cuboid, as shown in figure.

Hint: Use the fact that $a_1 = a_2$. This implies that the tensor has some symmetry, so that

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from the frame K you can reach the tensor in the K^\prime by rotating it with a matrix of the form:

$$R_{\phi} = \begin{bmatrix} \cos\phi & 0 & -\sin\phi \\ 0 & 1 & 0 \\ \sin\phi & 0 & \cos\phi \end{bmatrix}$$



Abbildung 1: The two frames K and K' mentioned above.