Exercise 4: Exact diagonalization

In this exercise, we consider the transverse field Ising model in 1D, given by the Hamiltonian

$$H = -J \sum_{j=0}^{L} \sigma_j^x \sigma_{j+1}^x - g \sum_{j=0}^{L-1} \sigma_j^z,$$
(1)

where we used periodic boundary conditions, $\sigma_L^x \equiv \sigma_0^x$. As discussed in the lecture, this model shows (for $L \to \infty$) a quantum phase transition at g = 1. To see this transition, we need to find the groundstate while tuning g (keeping $J \equiv 1$ as the unit of energy). The goal of this exercise is to represent eq. (1) as a sparse matrix, scipy.sparse.csr_matrix and diagonalize it with (a variant of) the Lanczos algorithm provided in scipy.

a) Define the 2×2 matrices

$$Id = \mathbb{1} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \qquad Sx = \sigma^x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \qquad Sz = \sigma^z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \qquad (2)$$

as scipy.sparse.csr_matrix

b) When we write σ_i^z in eq. (1), what we mean is

$$\sigma_j^z \equiv \mathbb{1} \otimes \dots \otimes \mathbb{1} \otimes \sigma^z \otimes \mathbb{1} \otimes \dots \otimes \mathbb{1} \tag{3}$$

where the σ^z is at the j^{th} position, and similar for σ_j^x . The operator σ_j^z corresponds thus to a $2^L \times 2^L$ matrix. To implement the tensor product, you can use succesive calls to scipy.sparse.kron(), e.g. to generate $\mathbb{1} \otimes \mathbb{1} \otimes \sigma^z \otimes \mathbb{1}$:

```
full = scipy.sparse.kron(Id, Id, format='csr') # 1 1
full = scipy.spares.kron(full, Sz, format='csr') # 1 1 Sz
full = scipy.spares.kron(full, Id, format='csr') # 1 1 Sz 1
```

Write a function which returns (for given L) a list, which contains a representation of σ_j^z (in the form of a csr_matrix) as j^{th} entry of the list.

- c) Write a similar function returning the σ_j^x operators.
- d) Write a function gen_hamiltonian(sx_list, sz_list, g, J) generating the Hamiltonian as a csr_matrix of shape $2^L \times 2^L$, where the arguments sx_list, sz_list should be the lists generated by the functions defined in b) and c).

Hint: Addition and scalar multiplication of sparse matrices work as expected with + and *. For two csr_matrix, the code A*B gives the **matrix** product of A and B. (This is different from numpy arrays where A*B gives the element-wise product!) Also note $A \otimes B = (A \otimes 1)(1 \otimes B)$, hence $\sigma_j^x \sigma_{j+1}^x$ can be obtained by a matrix product of σ_j^x (in the sense of eq. (3)) with σ_{j+1}^x .

Check that the function works as expected, e.g., for L = 2 and g = 0.1, you should obtain

$$H(L=2,g=0.1,J=1) = \begin{pmatrix} -0.2 & 0 & 0 & -2\\ 0 & 0 & -2 & 0\\ 0 & -2 & 0 & 0\\ -2 & 0 & 0 & 0.2 \end{pmatrix}.$$
 (4)

Hint: You can convert the sparse H to a non-sparse numpy array with H.toarray() for the comparison. Remember to use periodic boundary conditions.

e) When you have *H* as a csr_matrix, you can use the function scipy.sparse.linalg.eigsh to obtain the ground state. This function uses (an improved version of) the Lanczos method discussed in the lecture. Read the documentation to find out how to obtain the 3 eigen states (and eigen values) with smallest (algebraic) eigen values.

Optionally: Compare the run time of scipy.sparse.linalg.eigsh with the full diagonalization performed by np.linalg.eigh (acting on a non-sparse numpy array) for L = 12.

f) For system sizes $L \in \{6, 8, 10, 12\}$, calculate the ground state for ≈ 20 values of $g \in [0, 2]$. Calculate and plot the largest-distance spin-spin correlation $C = \langle \psi_0 | \sigma_0^x \sigma_{L/2}^x | \psi_0 \rangle$ against g for the various system sizes.

Hint: To calculate C, apply the operators (= sparse matrices) σ_0^x and $\sigma_{L/2}^x$ to the ground state using matrix-vector products (implemented as op*v0) and calculate the overlap of the result with the the ground state (e.g., using np.inner).

g) Plot the excitation energies for the first two excited states versus g. Does the result coincide with your expectations?